

Problem 4.59

Suppose two spin-1/2 particles are known to be in the singlet configuration (Equation 4.176). Let $S_a^{(1)}$ be the component of the spin angular momentum of particle number 1 in the direction defined by the vector \mathbf{a} . Similarly, let $S_b^{(2)}$ be the component of 2's angular momentum in the direction \mathbf{b} . Show that

$$\langle S_a^{(1)} S_b^{(2)} \rangle = -\frac{\hbar^2}{4} \cos \theta, \quad (4.225)$$

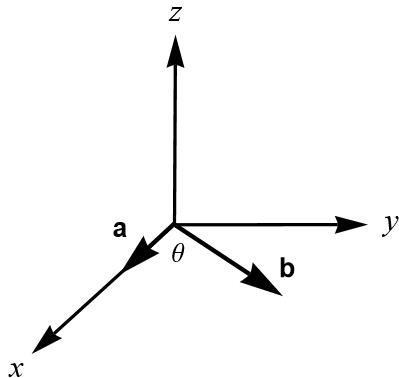
where θ is the angle between \mathbf{a} and \mathbf{b} .

Solution

Equation 4.225 has to hold regardless of which coordinate system is chosen.

Method 1

Choose a coordinate system where \mathbf{a} lies along the x -axis and \mathbf{b} lies in the xy -plane.



Evaluate the expectation value of $S_a^{(1)} S_b^{(2)}$ for the spin singlet, $|00\rangle$. Let δ_a and δ_b be the unit vectors in the directions of \mathbf{a} and \mathbf{b} , respectively.

$$\begin{aligned} \langle S_a^{(1)} S_b^{(2)} \rangle &= \langle 00 | S_a^{(1)} S_b^{(2)} | 00 \rangle \\ &= \langle 00 | (\delta_a \cdot \mathbf{S}^{(1)}) (\delta_b \cdot \mathbf{S}^{(2)}) | 00 \rangle \\ &= \langle 00 | \left[(\cos 0) S_x^{(1)} + \left(\cos \frac{\pi}{2} \right) S_y^{(1)} + \left(\cos \frac{\pi}{2} \right) S_z^{(1)} \right] \\ &\quad \left[(\cos \theta) S_x^{(2)} + \cos \left(\frac{\pi}{2} - \theta \right) S_y^{(2)} + \left(\cos \frac{\pi}{2} \right) S_z^{(2)} \right] | 00 \rangle \\ &= \langle 00 | S_x^{(1)} \left[(\cos \theta) S_x^{(2)} + (\sin \theta) S_y^{(2)} \right] | 00 \rangle \\ &= \langle 00 | S_x^{(1)} \left[(\cos \theta) S_x^{(2)} + (\sin \theta) S_y^{(2)} \right] \frac{1}{\sqrt{2}} (| \uparrow\downarrow \rangle - | \downarrow\uparrow \rangle) \end{aligned}$$

Write the spin operators in terms of the raising and lowering operators: $S_x = \frac{1}{2}(S_+ + S_-)$ and $S_y = \frac{1}{2i}(S_+ - S_-)$.

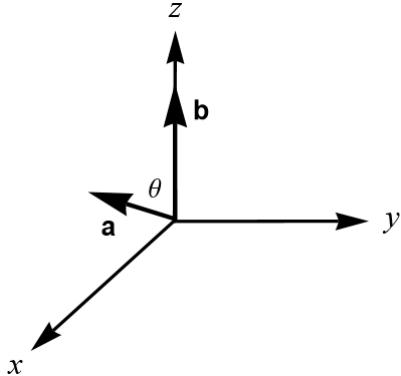
$$\begin{aligned}
\langle S_a^{(1)} S_b^{(2)} \rangle &= \left\langle 00 \left| \frac{1}{2}(S_+^{(1)} + S_-^{(1)}) \left[(\cos \theta) \frac{1}{2}(S_+^{(2)} + S_-^{(2)}) + (\sin \theta) \frac{1}{2i}(S_+^{(2)} - S_-^{(2)}) \right] \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \right. \right. \\
&= \left\langle 00 \left| (S_+^{(1)} + S_-^{(1)}) \left[(\cos \theta - i \sin \theta) S_+^{(2)} + (\cos \theta + i \sin \theta) S_-^{(2)} \right] \frac{1}{4\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \right. \right. \\
&= \left\langle 00 \left| (S_+^{(1)} + S_-^{(1)}) (e^{-i\theta} S_+^{(2)} + e^{i\theta} S_-^{(2)}) \frac{1}{4\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \right. \right. \\
&= \frac{1}{4\sqrt{2}} \left\langle 00 \left| (S_+^{(1)} + S_-^{(1)}) \left[e^{-i\theta} S_+^{(2)} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) + e^{i\theta} S_-^{(2)} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \right] \right. \right. \\
&= \frac{1}{4\sqrt{2}} \left\langle 00 \left| (S_+^{(1)} + S_-^{(1)}) \left[e^{-i\theta} |\uparrow\rangle (S_+^{(2)} |\downarrow\rangle) - e^{-i\theta} |\downarrow\rangle (S_+^{(2)} |\uparrow\rangle) \right. \right. \right. \\
&\quad \left. \left. \left. + e^{i\theta} |\uparrow\rangle (S_-^{(2)} |\downarrow\rangle) - e^{i\theta} |\downarrow\rangle (S_-^{(2)} |\uparrow\rangle) \right] \right. \right. \\
&= \frac{1}{4\sqrt{2}} \left\langle 00 \left| (S_+^{(1)} + S_-^{(1)}) \left[e^{-i\theta} |\uparrow\rangle (\hbar |\uparrow\rangle) - e^{-i\theta} |\downarrow\rangle (\mathbf{0}) + e^{i\theta} |\uparrow\rangle (\mathbf{0}) - e^{i\theta} |\downarrow\rangle (\hbar |\downarrow\rangle) \right] \right. \right. \\
&= \frac{\hbar}{4\sqrt{2}} \left\langle 00 \left| (S_+^{(1)} + S_-^{(1)}) \left(e^{-i\theta} |\uparrow\uparrow\rangle - e^{i\theta} |\downarrow\downarrow\rangle \right) \right. \right. \\
&= \frac{\hbar}{4\sqrt{2}} \left\langle 00 \left| \left[\left(e^{-i\theta} S_+^{(1)} |\uparrow\uparrow\rangle - e^{i\theta} S_+^{(1)} |\downarrow\downarrow\rangle \right) + \left(e^{-i\theta} S_-^{(1)} |\uparrow\uparrow\rangle - e^{i\theta} S_-^{(1)} |\downarrow\downarrow\rangle \right) \right] \right. \right. \\
&= \frac{\hbar}{4\sqrt{2}} \left\langle 00 \left| \left[e^{-i\theta} (S_+^{(1)} |\uparrow\rangle) |\uparrow\rangle - e^{i\theta} (S_+^{(1)} |\downarrow\rangle) |\downarrow\rangle + e^{-i\theta} (S_-^{(1)} |\uparrow\rangle) |\uparrow\rangle - e^{i\theta} (S_-^{(1)} |\downarrow\rangle) |\downarrow\rangle \right] \right. \right. \\
&= \frac{\hbar}{4\sqrt{2}} \left\langle 00 \left| \left[e^{-i\theta} (\mathbf{0}) |\uparrow\rangle - e^{i\theta} (\hbar |\uparrow\rangle) |\downarrow\rangle + e^{-i\theta} (\hbar |\downarrow\rangle) |\uparrow\rangle - e^{i\theta} (\mathbf{0}) |\downarrow\rangle \right] \right. \right. \\
&= \frac{\hbar^2}{4\sqrt{2}} \left\langle 00 \left| \left(-e^{i\theta} |\uparrow\downarrow\rangle + e^{-i\theta} |\downarrow\uparrow\rangle \right) \right. \right. \\
&= \frac{\hbar^2}{4\sqrt{2}} \langle 00 | \left[(\cos \theta - i \sin \theta) |\downarrow\uparrow\rangle - (\cos \theta + i \sin \theta) |\uparrow\downarrow\rangle \right] \\
&= \frac{\hbar^2}{4} \langle 00 | \left[\frac{\cos \theta}{\sqrt{2}} (|\downarrow\uparrow\rangle - |\uparrow\downarrow\rangle) - \frac{i \sin \theta}{\sqrt{2}} (|\downarrow\uparrow\rangle + |\uparrow\downarrow\rangle) \right] \\
&= \frac{\hbar^2}{4} \langle 00 | \left[\cos \theta (-|00\rangle) - i \sin \theta (|10\rangle) \right]
\end{aligned}$$

Therefore, because the states are orthonormal,

$$\begin{aligned}\left\langle S_a^{(1)} S_b^{(2)} \right\rangle &= -\frac{\hbar^2}{4} \cos \theta \langle 0 0 | 0 0 \rangle - \frac{i\hbar^2}{4} \sin \theta \langle 0 0 | 1 0 \rangle \\ &= -\frac{\hbar^2}{4} \cos \theta (1) - \frac{i\hbar^2}{4} \sin \theta (0) \\ &= -\frac{\hbar^2}{4} \cos \theta.\end{aligned}$$

Method 2

Choose a coordinate system where \mathbf{b} lies along the z -axis and \mathbf{a} lies in the xz -plane.



Evaluate the expectation value of $S_a^{(1)} S_b^{(2)}$ for the spin singlet, $|00\rangle$. Let δ_a and δ_b be the unit vectors in the directions of \mathbf{a} and \mathbf{b} , respectively.

$$\begin{aligned}
 \langle S_a^{(1)} S_b^{(2)} \rangle &= \langle 00 | S_a^{(1)} S_b^{(2)} | 00 \rangle \\
 &= \langle 00 | (\delta_a \cdot \mathbf{S}^{(1)}) (\delta_b \cdot \mathbf{S}^{(2)}) | 00 \rangle \\
 &= \langle 00 | \left[\cos\left(\frac{\pi}{2} - \theta\right) S_x^{(1)} + \left(\cos \frac{\pi}{2}\right) S_y^{(1)} + (\cos \theta) S_z^{(1)} \right] \\
 &\quad \left[\left(\cos \frac{\pi}{2}\right) S_x^{(2)} + \left(\cos \frac{\pi}{2}\right) S_y^{(2)} + (\cos 0) S_z^{(2)} \right] | 00 \rangle \\
 &= \langle 00 | \left[(\sin \theta) S_x^{(1)} + (\cos \theta) S_z^{(1)} \right] S_z^{(2)} | 00 \rangle \\
 &= \left\langle 00 \left| \left[(\sin \theta) S_x^{(1)} + (\cos \theta) S_z^{(1)} \right] S_z^{(2)} \right| \frac{1}{\sqrt{2}} (| \uparrow \downarrow \rangle - | \downarrow \uparrow \rangle) \right. \\
 &= \frac{1}{\sqrt{2}} \left\langle 00 \left| \left[(\sin \theta) S_x^{(1)} + (\cos \theta) S_z^{(1)} \right] \left[| \uparrow \rangle \left(S_z^{(2)} | \downarrow \rangle \right) - | \downarrow \rangle \left(S_z^{(2)} | \uparrow \rangle \right) \right] \right. \\
 &= \frac{1}{\sqrt{2}} \left\langle 00 \left| \left[(\sin \theta) S_x^{(1)} + (\cos \theta) S_z^{(1)} \right] \left[| \uparrow \rangle \left(-\frac{\hbar}{2} | \downarrow \rangle \right) - | \downarrow \rangle \left(\frac{\hbar}{2} | \uparrow \rangle \right) \right] \right. \\
 &= -\frac{\hbar}{2\sqrt{2}} \left\langle 00 \left| \left[(\sin \theta) S_x^{(1)} + (\cos \theta) S_z^{(1)} \right] (| \uparrow \downarrow \rangle + | \downarrow \uparrow \rangle) \right. \right. \\
 &= -\frac{\hbar}{2\sqrt{2}} \left\langle 00 \left| \left[(\sin \theta) \frac{1}{2} (S_+^{(1)} + S_-^{(1)}) (| \uparrow \downarrow \rangle + | \downarrow \uparrow \rangle) + (\cos \theta) S_z^{(1)} (| \uparrow \downarrow \rangle + | \downarrow \uparrow \rangle) \right] \right. \right.
 \end{aligned}$$

Therefore, because the states are orthonormal,

$$\begin{aligned}
 \langle S_a^{(1)} S_b^{(2)} \rangle &= -\frac{\hbar}{2\sqrt{2}} \left\langle 00 \left| \left[\frac{\sin \theta}{2} (S_+^{(1)} |\uparrow\downarrow\rangle + S_+^{(1)} |\downarrow\uparrow\rangle + S_-^{(1)} |\uparrow\downarrow\rangle + S_-^{(1)} |\downarrow\uparrow\rangle) \right. \right. \right. \\
 &\quad \left. \left. \left. + \cos \theta (S_z^{(1)} |\uparrow\downarrow\rangle + S_z^{(1)} |\downarrow\uparrow\rangle) \right] \right\} \\
 &= -\frac{\hbar}{2\sqrt{2}} \left\langle 00 \left| \left\{ \frac{\sin \theta}{2} \left[(S_+^{(1)} |\uparrow\rangle) |\downarrow\rangle + (S_+^{(1)} |\downarrow\rangle) |\uparrow\rangle + (S_-^{(1)} |\uparrow\rangle) |\downarrow\rangle + (S_-^{(1)} |\downarrow\rangle) |\uparrow\rangle \right] \right. \right. \\
 &\quad \left. \left. + \cos \theta \left[(S_z^{(1)} |\uparrow\rangle) |\downarrow\rangle + (S_z^{(1)} |\downarrow\rangle) |\uparrow\rangle \right] \right\} \right\} \\
 &= -\frac{\hbar}{2\sqrt{2}} \left\langle 00 \left| \left\{ \frac{\sin \theta}{2} \left[(\mathbf{0}) |\downarrow\rangle + (\hbar |\uparrow\rangle) |\uparrow\rangle + (\hbar |\downarrow\rangle) |\downarrow\rangle + (\mathbf{0}) |\uparrow\rangle \right] \right. \right. \\
 &\quad \left. \left. + \cos \theta \left[\left(\frac{\hbar}{2} |\uparrow\rangle \right) |\downarrow\rangle + \left(-\frac{\hbar}{2} |\downarrow\rangle \right) |\uparrow\rangle \right] \right\} \right\} \\
 &= -\frac{\hbar^2}{4} \left\langle 00 \left| \left[\frac{\sin \theta}{\sqrt{2}} (|\uparrow\uparrow\rangle + |\downarrow\downarrow\rangle) + \frac{\cos \theta}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \right] \right\} \\
 &= -\frac{\hbar^2}{4} \left\langle 00 \left| \left[\frac{\sin \theta}{\sqrt{2}} (|11\rangle + |1-1\rangle) + \cos \theta (|00\rangle) \right] \right\} \\
 &= -\frac{\hbar^2}{4} \left[\frac{\sin \theta}{\sqrt{2}} (\langle 00|11\rangle + \langle 00|1-1\rangle) + \cos \theta \langle 00|00\rangle \right] \\
 &= -\frac{\hbar^2}{4} \left[\frac{\sin \theta}{\sqrt{2}} (0 + 0) + \cos \theta (1) \right] \\
 &= -\frac{\hbar^2}{4} \cos \theta.
 \end{aligned}$$